

APPLICATION OF THE MONTE CARLO METHOD
TO THE PROBLEM OF FLOW IN THE
BOUNDARY LAYER OF A TWO-DIMENSIONAL
TURBULENT JET

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A precise solution is found to the problem of flow in the boundary layer of a two-dimensional turbulent jet. The method of [1] is used here.

The flow of an incompressible fluid in the boundary layer of a two-dimensional turbulent jet is described by the following equations and boundary conditions:

$$u_\beta u_{\alpha,\beta} + P_{,\alpha} = \tau_{\alpha\beta,\beta}; \quad (1)$$

$$u_{\sigma,\sigma} = 0; \quad (2)$$

$$\tau_{\alpha\beta} = -l_{\alpha\beta}^2 u_{1,2}^2; \quad (3)$$

$$\left. \begin{aligned} u_1|_{\Gamma_0} &= 1, \\ u_2|_{\Gamma_0} &= u_{1,2}|_{\Gamma_0} = u_1|_{\Gamma_1} = u_{1,2}|_{\Gamma_1} = 0 \end{aligned} \right\}. \quad (4)$$

The summation rule has been adopted here with respect to the repetitive Greek letter subscripts which assume the values 1, 2. The symbol after a comma denotes the derivative with respect to the given coordinate and u_α are the velocity components. On the basis of the similarity theory, we assume that u_α , P , and $x_1^{-1} l_{\alpha\beta}$ are functions of the ratio $\eta = x_1^{-1} x_2$. The flow function will be sought in the form

$$\psi = x_1 F(\eta).$$

Adopting the obvious rules of differentiation

$$A_{,1} = -A_{,\eta} \eta x_1^{-1}; \quad A_{,2} = A_{,\eta} x_1^{-1}$$

and considering that

$$\begin{aligned} u_1 &= \psi_{,2}, \quad u_2 = -\psi_{,1}, \\ l_{\alpha\beta} &= \mu_{\alpha\beta}(\eta) x_1, \end{aligned}$$

one can reduce systems (1)-(4) to a single third-order ordinary differential equation (a prime sign indicates a derivative with respect to η)

$$\begin{aligned} Q &= [(1-\eta^2) \mu_{12}^2 + \eta(\mu_{22}^2 - \mu_{11}^2)] F''' + [(1-\eta^2) \mu_{12} \mu_{12}' \\ &+ \eta(\mu_{22} \mu_{22}' - \mu_{11} \mu_{11}')] F'' - \frac{1}{2} (1+\eta^2) F = 0. \end{aligned} \quad (5)$$

The boundary conditions (4) become

$$F|_{\alpha=0} = \eta_0; \quad F'|_{\alpha=0} = 1; \quad F''|_{\alpha=0} = F'|_{\alpha=1} = F''|_{\alpha=1} = 0. \quad (6)$$

Here $\alpha = (\eta - \eta_0)/(\eta_1 - \eta_0)$; $\alpha \in [0; +1]$, while parameters η_0 and η_1 define the location of boundaries Γ_0 and Γ_1 . To conditions (6) we add two others:

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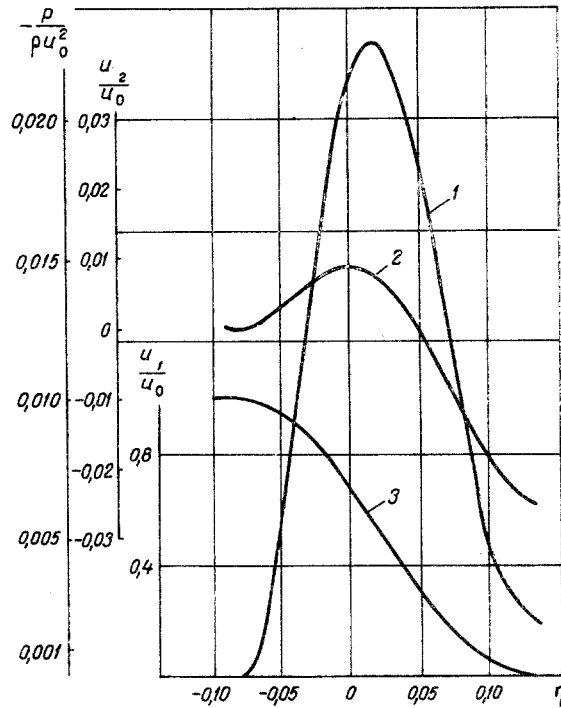


Fig. 1. Functions of $\eta = x_2/x_1$: 1) $-P/\rho u_0^2$; 2) u_2/u_0 ; 3) u_1/u_0 (given: $\mu_{11} = 0.030$, $\mu_{12} = \mu_{21} = 0.014$, $\mu_{22} = 0.020$, $\nu = 50$; obtained: $\eta_0 = -0.0776$, $\eta_1 = +0.1382$, $k_0 = -130.6$, $k_1 = +66.51$).

$$F''|_{\alpha=0} = k_0; F''|_{\alpha=1} = k_1,$$

where k_0, k_1 are constants not yet known.

The solution to the boundary-value problems (5)-(6) will be sought in the form of a polynomial

$$F = \varphi_{-1} + \varphi_0 \alpha + \frac{1}{2} \varphi_1 \alpha^2 + \frac{1}{3} \varphi_2 \alpha^3 + \frac{1}{4} \varphi_3 \alpha^4 + \frac{1}{5} \varphi_4 \alpha^5 + \frac{1}{6} \varphi_5 \alpha^6, \quad (7)$$

whose coefficients are defined as follows:

$$\varphi_{-1} = \eta_0, \varphi_0 = \Delta = \eta_1 - \eta_0, \varphi_1 = 0, \varphi_2 = \frac{1}{2} \Delta^3 k_0,$$

$$\varphi_3 = \left[-10 + (k_1 - 3k_0) \frac{\Delta^2}{2} \right] \Delta,$$

$$\varphi_4 = \left[15 + (3k_0 - 2k_1) \frac{\Delta^2}{2} \right] \Delta,$$

$$\varphi_5 = \left[-6 + (k_1 - k_0) \frac{\Delta^2}{2} \right] \Delta.$$

In this case the boundary conditions are satisfied automatically and the four parameters η_0, η_1, k_0, k_1 remain unconstrained.

In accordance with [1], the values of these parameters at which the quantity

$$\Phi = N^{-1} \sum_{s=1}^N Q_s^2 \quad (8)$$

becomes minimum will be considered optimal. Here N is the total number of control points on the interval $[\eta_0, \eta_1]$. Their coordinates were calculated by the formula

$$\eta^{(s)} = \eta_0 + sN^{-1}(\eta_1 - \eta_0),$$

with s denoting the consecutive number of a point and Q_s denoting the magnitude of the error incurred by inserting polynomial (7) into Eq. (5) at $\eta = \eta^{(s)}$. In accordance with [3], the following values were taken for $\mu_{\alpha\beta}$: $\mu_{11} = 0.03$, $\mu_{12} = 0.014$, $\mu_{22} = 0.02$. Equation (8) was minimized by the method of random tracking [2]. On the basis of optimal values for η_0 , η_1 , k_0 , and k_1 with the aid of expressions

$$u_1/u_0 = F'; \quad u_2/u_0 = \eta F' - F;$$

$$\frac{P}{\rho u_0^2} = 2 \int_{\eta_0}^{\eta_1} \left[\frac{1}{2} \eta F + \eta \mu_{12} (\mu'_{12} F'' + \mu_{12} F''') \right. \\ \left. - \mu_{22} (\mu'_{22} F'' + \mu_{22} F''') \right] F'' d\eta$$

curves have been plotted as shown in Fig. 1. They come close to the empirical curves in [3].

Evidently, a further refinement is possible by increasing the number of terms in expression (7) and by changing conditions (6) to more stringent ones (with transverse flow also taken into account).

NOTATION

x_1, x_2	are the Cartesian coordinates;
u_1, u_2	are the velocity components;
P	is the pressure, referred to a plane surface;
$\eta = x_2/x_1$	is the dimensionless coordinate;
Γ_0, Γ_1	are the inside and outside boundary of a layer;
η_0, η_1	are their respective coordinates;
$\alpha = (\eta - \eta_0)/(\eta_1 - \eta_0)$	is the coordinate referred to interval [0, +1];
$\psi = x_1 F(\eta)$	is the flow function;
$\tau_{\alpha\beta} = -\rho u_{1,2}^2 = -\mu_{\alpha\beta}^2(\eta) x_1^2 u_{1,2}^2$	are the components of the tensor of turbulent stresses, referred to the density;
F	is the power polynomial representing the solution;
$\varphi_{-1}, \varphi_0, \varphi_1, \dots, \varphi_5$	are the coefficients of the power polynomial;
Φ	is the functional to be minimized.

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